

# Positive and nodal single-layered solutions to supercritical elliptic problems above the higher critical exponents

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*To Jean Mawhin on his 75th birthday, with great appreciation*

ABSTRACT. *We study the problem*

$$-\Delta v + \lambda v = |v|^{p-2} v \text{ in } \Omega, \quad v = 0 \text{ on } \partial\Omega,$$

*for  $\lambda \in \mathbb{R}$  and supercritical exponents  $p$ , in domains of the form*

$$\Omega := \{(y, z) \in \mathbb{R}^{N-m-1} \times \mathbb{R}^{m+1} : (y, |z|) \in \Theta\},$$

*where  $m \geq 1$ ,  $N - m \geq 3$ , and  $\Theta$  is a bounded domain in  $\mathbb{R}^{N-m}$  whose closure is contained in  $\mathbb{R}^{N-m-1} \times (0, \infty)$ . Under some symmetry assumptions on  $\Theta$ , we show that this problem has infinitely many solutions for every  $\lambda$  in an interval which contains  $[0, \infty)$  and  $p > 2$  up to some number which is larger than the  $(m+1)^{\text{st}}$  critical exponent  $2_{N,m}^* := \frac{2(N-m)}{N-m-2}$ . We also exhibit domains with a shrinking hole, in which there are a positive and a nodal solution which concentrate on a sphere, developing a single layer that blows up at an  $m$ -dimensional sphere contained in the boundary of  $\Omega$ , as the hole shrinks and  $p \rightarrow 2_{N,m}^*$  from above. The limit profile of the positive solution, in the transversal direction to the sphere of concentration, is a rescaling of the standard bubble, whereas that of the nodal solution is a rescaling of a nonradial sign-changing solution to the problem*

$$-\Delta u = |u|^{2_n^*-2} u, \quad u \in D^{1,2}(\mathbb{R}^n),$$

*where  $2_n^* := \frac{2n}{n-2}$  is the critical exponent in dimension  $n$ .*

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## 1. Introduction

We study the existence and concentration behavior of solutions to the problem

$$\begin{cases} -\Delta v + \lambda v = |v|^{p-2} v & \text{in } \Omega, \\ v = 0 & \text{on } \partial\Omega, \end{cases} \quad (\wp_p)$$

where  $\Omega$  is a bounded domain in  $\mathbb{R}^N$ ,  $\lambda \in \mathbb{R}$ , and  $p$  is supercritical, i.e., it is larger than the critical Sobolev exponent  $2_N^* := \frac{2N}{N-2}$  for  $N \geq 3$ . We shall consider domains of the form

$$\Omega := \{(y, z) \in \mathbb{R}^{N-m-1} \times \mathbb{R}^{m+1} : (y, |z|) \in \Theta\}, \quad (1)$$

where  $m \geq 1$ ,  $N - m \geq 3$ , and  $\Theta$  is a bounded domain in  $\mathbb{R}^{N-m}$  whose closure is contained in  $\mathbb{R}^{N-m-1} \times (0, \infty)$ .

In domains of this type, the true critical exponent is  $2_{N,m}^* := \frac{2(N-m)}{N-m-2}$ , which is the critical Sobolev exponent in the dimension of  $\Theta$  and is larger than  $2_N^*$ . Indeed, one can easily verify that the solutions to the problem  $(\wp_p)$  which are radial in the variable  $z$ , correspond to the solutions of the problem

$$\begin{cases} -\operatorname{div}(f(x)u) + \lambda f(x)u = f(x)|u|^{p-2}u & \text{in } \Theta, \\ u = 0 & \text{on } \partial\Theta, \end{cases} \quad (2)$$

where  $f(x_1, \dots, x_{N-m}) = x_{N-m}^m$ . Standard variational methods show that this last problem has infinitely many solutions for  $p \in (2, 2_{N-m}^*)$ , hence, also does the problem  $(\wp_p)$ . On the other hand, Passaseo showed in [18, 19] that, if  $\lambda = 0$  and  $\Theta$  is a ball centered on the half-line  $\{0\} \times (0, \infty)$ , then the problem  $(\wp_p)$  does not have a nontrivial solution for  $p \geq 2_{N-m}^* = 2_{N,m}^*$ . The number  $2_{N,m}^*$  has been called the  $(m+1)^{\text{st}}$  *critical exponent* in dimension  $N$ .

The concentration behavior of solutions to the problem  $(\wp_p)$  for  $\lambda = 0$  and  $p \in (2, 2_{N,m}^*)$ , as  $p \rightarrow 2_{N,m}^*$  from below, has been investigated in several papers. In [10], del Pino, Musso and Pacard exhibited positive solutions which concentrate and blow up at a nondegenerate closed geodesic in  $\partial\Omega$ , as  $p$  approaches the second critical exponent  $2_{N,1}^*$  from below. For any  $m \geq 1$ , positive and sign-changing solutions in domains of the form (1) were constructed in [1, 13]. These solutions concentrate and blow up at one or several  $m$ -dimensional spheres, as  $p \rightarrow 2_{N,m}^*$  from below. In all of these cases the limit profile of the solutions, in the transversal direction to each sphere of concentration, is a sum of rescalings of  $\pm U$ , where

$$U(x) := [n(n-2)]^{(n-2)/4} \left( \frac{1}{1+|x|^2} \right)^{(n-2)/2}$$

is the standard bubble in dimension  $n := N - m$ , which is the only positive solution to the limit problem

$$-\Delta u = |u|^{2_n^*-2} u, \quad u \in D^{1,2}(\mathbb{R}^n), \quad (3)$$

up to translation and dilation.

It was recently shown in [4] that there exist nonradial sign-changing solutions to the problem (3), that do not resemble a sum of rescaled positive and negative standard bubbles, which occur as limit profiles for concentration of sign-changing solutions to the problem  $(\varphi_p)$  that blow up at a single point, as  $p \rightarrow 2_N^*$  from below. For the higher critical exponents  $2_{N,m}^*$  with  $m \geq 1$ , it was shown in [5] that for every  $\lambda$  in some interval which contains  $[0, \infty)$  there are sign-changing solutions to the problem  $(\varphi_p)$ , in domains of the form (1), which concentrate and blow up at an  $m$ -dimensional sphere, as  $p \rightarrow 2_{N,m}^*$  from below, whose limit profile in the transversal direction to the sphere of concentration is a nonradial sign-changing solution to (3), like those found in [4].

The study of concentration phenomena for  $p$  approaching  $2_N^*$  from above, is a much more delicate issue, beginning with the fact that solutions to  $(\varphi_p)$  for  $p > 2_N^*$  do not always exist. For  $\lambda = 0$ , standard bubbles were used as basic cells in [8, 9, 16, 20] to construct positive solutions to the slightly supercritical problem  $(\varphi_p)$  with  $p = 2_N^* + \varepsilon$ , for small enough  $\varepsilon > 0$ , in domains with a hole, using the Lyapunov-Schmidt reduction method. These solutions blow up, as  $\varepsilon \rightarrow 0$ , and their limit profile at each blow-up point is a rescaling of the standard bubble. Solutions in some contractible domains were constructed in [14, 15].

Quite recently, sign-changing solutions to the slightly supercritical problem  $(\varphi_p)$  with  $p = 2_N^* + \varepsilon$ ,  $\varepsilon > 0$ , were exhibited by Musso and Wei [17] in domains with a small fixed hole, and by Clapp and Pacella [6] in domains with a shrinking hole. The solutions obtained in [17] concentrate at two different points in the domain, as  $\varepsilon \rightarrow 0$ , and their limit profile at each of them is a rescaling of one of the sign-changing solutions to the limit problem (3) in  $\mathbb{R}^N$  constructed by del Pino, Musso, Pacard and Pistoia in [11], which resemble a large number of negative bubbles, placed evenly along a circle, surrounding a positive bubble, placed at its center. On the other hand, the sign-changing solutions exhibited in [6] concentrate at a single point in the interior of the shrinking hole, as the hole shrinks and  $\varepsilon \rightarrow 0$ , and their limit profile is a rescaling of a nonradial sign-changing solution to (3) like those found in [4].

For  $m \geq 1$ , the existence of solutions for  $p = 2_{N,m}^* + \varepsilon$  and their concentration behavior seems to be, so far, an open question; see Problem 4 in [7]. In this paper we will show that, under some symmetry assumptions, the problem  $(\varphi_p)$  has infinitely many solutions in domains of the form (1) for  $p > 2_{N,m}^*$ , up to some value which depends on the symmetries; see Theorem 2.3. We will also exhibit domains with a shrinking hole, in which there are positive and

sign-changing solutions which concentrate and blow up at an  $m$ -dimensional sphere contained in the boundary of  $\Omega$ , as the hole shrinks and  $p \rightarrow 2_{N,m}^*$  from above. The limit profile of the positive solutions, in the direction transversal to the sphere of concentration, will be a rescaling of the standard bubble, whereas that of the sign-changing ones will resemble one of the solutions to (3) that were found in [4].

We give, next, some examples of our results. For  $n := N - m$ , let  $B$  be an  $n$ -dimensional ball of radius  $\delta_0$ , centered on the half-line  $\{0\} \times (0, \infty)$ , whose closure is contained in the half-space  $\mathbb{R}^{n-1} \times (0, \infty)$ . We write the points in  $\mathbb{R}^{n-1} \times (0, \infty)$  as  $(y, t)$  with  $y \in \mathbb{R}^{n-1}$ ,  $t \in (0, \infty)$  and we set

$$\begin{aligned} B_\delta &:= \{(y, t) \in B : |y| > \delta\} \quad \text{if } \delta \in (0, \delta_0), & B_0 &:= B, \\ \Omega_\delta &:= \{(y, z) \in \mathbb{R}^{n-1} \times \mathbb{R}^{m+1} : (y, |z|) \in B_\delta\}, & \Omega &:= \Omega_0. \end{aligned}$$

We denote by  $O(k)$  the group of all linear isometries of  $\mathbb{R}^k$  and, for  $v \in D^{1,2}(\mathbb{R}^N)$ , we write

$$\|v\| := \left( \int_{\mathbb{R}^N} |\nabla v|^2 \right)^{1/2}.$$

The following results establish the existence of positive and sign-changing solutions to the problem  $(\wp_p)$  in  $\Omega_\delta$  and describe their limit profile as  $\delta \rightarrow 0$  and  $p \rightarrow 2_{N,m}^*$  from above. They are special cases of Theorems 2.3 and 4.4, which apply to more general domains, and are stated and proved in Sections 2 and 4, respectively.

**THEOREM 1.1.** *There exists  $\lambda_* \leq 0$  such that, for each  $\lambda \in (\lambda_*, \infty) \cup \{0\}$ ,  $\delta \in (0, \delta_0)$  and  $p \in (2, \infty)$ , the problem  $(\wp_p)$  has a positive solution  $v_{\delta,p}$  in  $\Omega_\delta$  which satisfies*

$$v_{\delta,p}(\gamma y, \varrho z) = v_{\delta,p}(y, z) \quad \forall \gamma \in O(n-1), \varrho \in O(m+1), (y, z) \in \Omega_\delta,$$

and has minimal energy among all nontrivial solutions to  $(\wp_p)$  in  $\Omega_\delta$  with these symmetries.

Moreover, there exist sequences  $(\delta_k)$  in  $(0, \delta_0)$ ,  $(p_k)$  in  $(2_{N,m}^*, \infty)$ ,  $(\varepsilon_k)$  in  $(0, \infty)$  and  $(\zeta_k)$  in  $B \cap [\{0\} \times (0, \infty)]$  such that

(i)  $\delta_k \rightarrow 0$ ,  $p_k \rightarrow 2_{N,m}^*$ ,  $\varepsilon_k^{-1} \text{dist}(\zeta_k, \partial\Theta) \rightarrow \infty$ , and  $\zeta_k \rightarrow \zeta$  with

$$\text{dist}(\zeta, \mathbb{R}^{n-1} \times \{0\}) = \text{dist}(B, \mathbb{R}^{n-1} \times \{0\}),$$

(ii)  $\lim_{k \rightarrow \infty} \left\| v_{\delta_k, p_k} - \tilde{U}_{\varepsilon_k, \zeta_k} \right\| = 0$ , where

$$\tilde{U}_{\varepsilon_k, \zeta_k}(y, z) := \varepsilon_k^{(2-n)/2} U \left( \frac{(y, |z|) - \zeta_k}{\varepsilon_k} \right)$$

and  $U$  is the standard bubble in dimension  $n$ .

The number  $\lambda_*$  is negative if  $m \geq 2$ .

The solutions given by the previous theorem concentrate on an  $m$ -dimensional sphere, developing a positive layer which blows up at an  $m$ -dimensional sphere contained in the boundary of  $\Omega$  and located at minimal distance to the plane of rotation  $\mathbb{R}^{n-1} \times \{0\}$ . The asymptotic profile of each layer in the transversal direction to its sphere of concentration is a rescaling of the standard bubble.

The next theorem gives sign-changing solutions to the problem  $(\wp_p)$  with a different type of asymptotic profile. For  $n \geq 5$ , we write  $\mathbb{R}^{n-1} \equiv \mathbb{C}^2 \times \mathbb{R}^{n-5}$ , and the points in  $\mathbb{R}^{n-1}$  as  $y = (\eta, \xi)$ , with  $\eta = (\eta_1, \eta_2) \in \mathbb{C}^2$ ,  $\xi \in \mathbb{R}^{n-5}$ .

**THEOREM 1.2.** *Assume that  $n = 5$  or  $n \geq 7$ . Then, there exists  $\lambda_* \leq 0$  such that, for each  $\lambda \in (\lambda_*, \infty) \cup \{0\}$ ,  $\delta \in (0, \delta_0)$  and  $p \in (2, 2_{N, m+1}^*)$ , the problem  $(\wp_p)$  has a nontrivial sign-changing solution  $w_{\delta, p}$  in  $\Omega_\delta$  which satisfies*

$$w_{\delta, p}(\eta, \xi, z) = w_{\delta, p}(e^{i\vartheta}\eta, \alpha\xi, \varrho z), \quad w_{\delta, p}(\eta_1, \eta_2, \xi, z) = -w_{\delta, p}(-\bar{\eta}_2, \bar{\eta}_1, \xi, z),$$

for every  $\vartheta \in [0, \pi)$ ,  $\alpha \in O(n-5)$ ,  $\varrho \in O(m+1)$  and  $(y, z) \in \Omega_\delta$ , and which has minimal energy among all nontrivial solutions to  $(\wp_p)$  in  $\Omega_\delta$  with these symmetry properties.

Moreover, there exist sequences  $(\delta_k)$  in  $(0, \delta_0)$ ,  $(p_k)$  in  $(2_{N, m}^*, 2_{N, m+1}^*)$ ,  $(\varepsilon_k)$  in  $(0, \infty)$  and  $(\zeta_k)$  in  $B \cap [\{0\} \times (0, \infty)]$ , and a nontrivial sign-changing solution  $W$  to the limit problem (3), such that

$$(i) \quad \delta_k \rightarrow 0, \quad p_k \rightarrow 2_{N, m}^*, \quad \varepsilon_k^{-1} \text{dist}(\zeta_k, \partial\Theta) \rightarrow \infty, \quad \text{and} \quad \zeta_k \rightarrow \zeta \quad \text{with}$$

$$\text{dist}(\zeta, \mathbb{R}^{n-1} \times \{0\}) = \text{dist}(B, \mathbb{R}^{n-1} \times \{0\}),$$

$$(ii) \quad W(\eta, \xi, t) = W(e^{i\vartheta}\eta, \alpha\xi, t) \quad \text{and} \quad W(\eta_1, \eta_2, \xi, t) = -W(-\bar{\eta}_2, \bar{\eta}_1, \xi, t) \quad \text{for}$$

every  $\vartheta \in [0, \pi)$ ,  $\alpha \in O(n-5)$  and  $(y, t) \in \mathbb{R}^{n-1} \times \mathbb{R} \equiv \mathbb{R}^n$ , and  $W$  has minimal energy among all nontrivial solutions to (3) with these symmetry properties,

$$(iii) \quad \lim_{k \rightarrow \infty} \left\| w_{\delta_k, p_k} - \widetilde{W}_{\varepsilon_k, \zeta_k} \right\| = 0, \quad \text{where}$$

$$\widetilde{W}_{\varepsilon_k, \zeta_k}(y, z) := \varepsilon_k^{(2-n)/2} W \left( \frac{(y, |z|) - \zeta_k}{\varepsilon_k} \right).$$

The number  $\lambda_*$  is negative if  $m \geq 2$ .

The solutions given by the previous theorem concentrate on an  $m$ -dimensional sphere, developing a sign-changing layer which blows up at an  $m$ -dimensional sphere contained in the boundary of  $\Omega$  and located at minimal distance

to the plane of rotation  $\mathbb{R}^{n-1} \times \{0\}$ . The asymptotic profile of each layer in the transversal direction to its sphere of concentration is a rescaling of a nonradial sign-changing solution to the limit problem (3), like those found in [4].

As we mentioned before, the solutions to the anisotropic problem (2) give rise to solutions of the problem  $(\wp_p)$  in domains of the form (1). In Section 2 we will study a general anisotropic problem in an  $n$ -dimensional domain  $\Theta$ . We will assume that  $\Theta$  has some symmetries and we will establish the existence of infinitely many positive and sign-changing solutions to the anisotropic problem for supercritical exponents  $p > 2_n^*$ , up to some value which depends on the symmetries. These results extend those obtained in [6] for the problem with constant coefficients. In Section 3 we will describe the behavior of the minimizing sequences for the variational functional associated to the anisotropic problem for  $p = 2_n^*$ . These sequences, either converge to a solution, or they blow up. We will provide information on the location of the blow-up points and on the symmetries of the solutions to the limit problem (3) which occur as limit profiles. This will be used in Section 4 to obtain information on the concentration behavior and the limit profile of positive and sign-changing solutions to the problem  $(\wp_p)$  in domains with a shrinking hole, as the hole shrinks and  $p \rightarrow 2_{N,m}^*$  from above.

## 2. Symmetries and existence for supercritical problems

Let  $\Gamma$  be a closed subgroup of  $O(n)$  and  $\phi : \Gamma \rightarrow \mathbb{Z}_2$  be a continuous homomorphism of groups. A function  $u : \mathbb{R}^n \rightarrow \mathbb{R}$  is said to be  $\phi$ -equivariant if

$$u(\gamma x) = \phi(\gamma)u(x) \quad \forall \gamma \in \Gamma, x \in \mathbb{R}^n. \quad (4)$$

If  $\phi$  is the trivial homomorphism, then (4) simply says that  $u$  is a  $\Gamma$ -invariant function, whereas, if  $\phi$  is surjective and  $u$  is not trivial, then (4) implies that  $u$  is sign-changing, nonradial and  $G$ -invariant, where  $G := \ker \phi$ .

Let  $\Theta$  be a bounded  $\Gamma$ -invariant domain in  $\mathbb{R}^n$ ,  $n \geq 3$ , and  $a \in C^1(\overline{\Theta})$ ,  $b, c \in C^0(\overline{\Theta})$  be  $\Gamma$ -invariant functions satisfying  $a, c > 0$  on  $\overline{\Theta}$ . We assume that

$$\text{there exists } x_0 \in \Theta \text{ such that } \{\gamma \in \Gamma : \gamma x_0 = x_0\} \subset \ker \phi. \quad (5)$$

This assumption guarantees that the space

$$D_0^{1,2}(\Theta)^\phi := \{u \in D_0^{1,2}(\Theta) : u \text{ is } \phi\text{-equivariant}\}$$

is infinite dimensional; see [3]. As usual,  $D_0^{1,2}(\Theta)$  denotes the closure of  $C_c^\infty(\Theta)$  in the Hilbert space

$$D^{1,2}(\mathbb{R}^n) := \{u \in L^{2_n^*}(\mathbb{R}^n) : \nabla u \in L^2(\mathbb{R}^n, \mathbb{R}^n)\}$$

equipped with the norm

$$\|u\| := \left( \int_{\Theta} |\nabla u|^2 \right)^{1/2}.$$

We shall also assume that the operator  $-\operatorname{div}(a\nabla) + b$  is coercive in  $D_0^{1,2}(\Theta)^\phi$ , i.e., that

$$\inf_{\substack{u \in D_0^{1,2}(\Theta)^\phi \\ u \neq 0}} \frac{\int_{\Theta} (a(x) |\nabla u|^2 + b(x) u^2) dx}{\int_{\Theta} |\nabla u|^2} > 0. \quad (6)$$

We set

$$\|u\|_{a,b}^2 := \int_{\Theta} (a(x) |\nabla u|^2 + b(x) u^2) dx, \quad |u|_{c;p}^p := \int_{\Theta} c(x) |u|^p dx.$$

Assumption (6) implies that  $\|\cdot\|_{a,b}$  is a norm in  $D_0^{1,2}(\Theta)^\phi$  which is equivalent to  $\|\cdot\|$ . Note that, as  $c > 0$ ,  $|\cdot|_{c;p}$  is equivalent to the standard norm in  $L^p(\Theta)$ , which we denote by  $|\cdot|_p$ .

Our aim is to establish the existence of solutions to the problem

$$\begin{cases} -\operatorname{div}(a(x)\nabla u) + b(x)u = c(x)|u|^{p-2}u & \text{in } \Theta, \\ u = 0 & \text{on } \partial\Theta. \\ u(\gamma x) = \phi(\gamma)u(x), & \forall \gamma \in \Gamma, \ x \in \Theta, \end{cases} \quad (7)$$

for every  $2 < p < 2_{n-d}^*$ , where

$$d := \min\{\dim(\Gamma x) : x \in \overline{\Theta}\},$$

$\Gamma x := \{\gamma x : \gamma \in \Gamma\}$  is the  $\Gamma$ -orbit of  $x$ ,  $2_k^* := \frac{2k}{k-2}$  if  $k \geq 3$  and  $2_k^* := \infty$  if  $k = 1, 2$ . Note that  $2_{n-d}^* > 2_n^*$  if  $d > 0$ .

A (weak) solution to the problem (7) is a function  $u \in D_0^{1,2}(\Theta)^\phi \cap L^p(\Theta)$  such that

$$\int_{\Theta} (a(x)\nabla u \cdot \nabla \psi + b(x)u\psi) dx - \int_{\Theta} c(x)|u|^{p-2}u\psi dx = 0 \quad \forall \psi \in \mathcal{C}_c^\infty(\Theta). \quad (8)$$

Proposition 2.1 of [6] asserts that  $D_0^{1,2}(\Theta)^\phi$  is continuously embedded in  $L^p(\Theta)$  for any real number  $p \in [1, 2_{n-d}^*]$ , and that the embedding is compact for  $p \in [1, 2_{n-d}^*)$ . The proof relies on a result by Hebey and Vaugon [12] which establishes these facts for  $\Gamma$ -invariant functions. Therefore, the functional

$$J_p(u) := \frac{1}{2} \|u\|_{a,b}^2 - \frac{1}{p} |u|_{c;p}^p$$

is well defined in the space  $D_0^{1,2}(\Theta)^\phi$  if  $p \in (2, 2_{n-d}^*]$ .

LEMMA 2.1. *For any real number  $p \in (2, 2_{n-d}^*]$  the critical points of the functional  $J_p$  in the space  $D_0^{1,2}(\Theta)^\phi$  are the solutions to the problem (7).*

*Proof.* Let  $u \in D_0^{1,2}(\Theta)^\phi$  be a critical point of  $J_p$  in  $D_0^{1,2}(\Theta)^\phi$ . Then,

$$J'_p(u)\vartheta = \int_{\Theta} (a(x)\nabla u \cdot \nabla \vartheta + b(x)u\vartheta - c(x)|u|^{p-2}u\vartheta) dx = 0 \quad \forall \vartheta \in D_0^{1,2}(\Theta)^\phi.$$

As  $D_0^{1,2}(\Theta)^\phi \subset L^p(\Theta)$  for  $1 \leq p \leq 2_{n-d}^*$  we need only to prove that  $u$  satisfies (8). Let  $\psi \in \mathcal{C}_c^\infty(\Theta)$ , and define

$$\tilde{\psi}(x) := \frac{1}{\mu(\Gamma)} \int_{\Gamma} \phi(\gamma)\psi(\gamma x) d\mu,$$

where  $\mu$  is the Haar measure on  $\Gamma$ . Note that  $\tilde{\psi} \in D_0^{1,2}(\Theta)^\phi$ . Observe also that, as  $u$  is  $\phi$ -equivariant, we have that

$$\phi(\gamma)\nabla u(x) = \nabla(u \circ \gamma)(x) = \gamma^{-1}\nabla u(\gamma x) \quad \forall \gamma \in \Gamma, x \in \Theta.$$

Since  $J'_p(u)\tilde{\psi} = 0$ , and  $a, b, c$  are  $\Gamma$ -invariant, using Fubini's theorem and performing a change of variable, we get

$$\begin{aligned} 0 &= \int_{\Theta} (a(x)\nabla u(x) \cdot \nabla \tilde{\psi}(x) + b(x)u(x)\tilde{\psi}(x) - c(x)|u(x)|^{p-2}u(x)\tilde{\psi}(x)) dx \\ &= \frac{1}{\mu(\Gamma)} \int_{\Theta} \int_{\Gamma} [a(x)\phi(\gamma)\nabla u(x) \cdot \gamma^{-1}\nabla \psi(\gamma x) + b(x)\phi(\gamma)u(x)\psi(\gamma x) \\ &\quad - c(x)|\phi(\gamma)u(x)|^{p-2}\phi(\gamma)u(x)\psi(\gamma x)] d\mu dx \\ &= \frac{1}{\mu(\Gamma)} \int_{\Gamma} \int_{\Theta} [a(x)\gamma^{-1}\nabla u(\gamma x) \cdot \gamma^{-1}\nabla \psi(\gamma x) + b(x)u(\gamma x)\psi(\gamma x) \\ &\quad - c(x)|u(\gamma x)|^{p-2}u(\gamma x)\psi(\gamma x)] dx d\mu \\ &= \frac{1}{\mu(\Gamma)} \int_{\Gamma} \int_{\Theta} [a(\gamma x)\nabla u(\gamma x) \cdot \nabla \psi(\gamma x) + b(\gamma x)u(\gamma x)\psi(\gamma x) \\ &\quad - c(\gamma x)|u(\gamma x)|^{p-2}u(\gamma x)\psi(\gamma x)] dx d\mu \\ &= \frac{1}{\mu(\Gamma)} \int_{\Gamma} d\mu \int_{\Theta} [a(\xi)\nabla u(\xi) \cdot \nabla \psi(\xi) + b(\xi)u(\xi)\psi(\xi) \\ &\quad - c(\xi)|u(\xi)|^{p-2}u(\xi)\psi(\xi)] d\xi \\ &= \int_{\Theta} [a(\xi)\nabla u(\xi) \cdot \nabla \psi(\xi) + b(\xi)u(\xi)\psi(\xi) - c(\xi)|u(\xi)|^{p-2}u(\xi)\psi(\xi)] d\xi. \end{aligned}$$

Therefore  $u$  is a solution to the problem (7).  $\square$



The nontrivial critical points of the functional  $J_p : D_0^{1,2}(\Theta)^\phi \rightarrow \mathbb{R}$  lie on the Nehari manifold

$$\mathcal{N}_p^\phi := \left\{ u \in D_0^{1,2}(\Theta)^\phi : \|u\|_{a,b}^2 = |u|_{c;p}^p, u \neq 0 \right\},$$

which is a  $\mathcal{C}^2$ -Hilbert manifold, radially diffeomorphic to the unit sphere in  $D_0^{1,2}(\Theta)^\phi$ , and a natural constraint for this functional. Set

$$\ell_p^\phi := \inf\{J_p(u) : u \in \mathcal{N}_p^\phi\}.$$

Then,  $\ell_p^\phi > 0$ . A *least energy solution* to the problem (7) is a minimizer for  $J_p$  on  $\mathcal{N}_p^\phi$ . The following result extends Theorem 2.3 in [6].

**THEOREM 2.2.** *If  $p \in (2, 2_{n-d}^*)$  then the problem (7) has a least energy solution, and an unbounded sequence of solutions.*

*Proof.* By Lemma 2.1, the critical points of the functional  $J_p$  in the space  $D_0^{1,2}(\Theta)^\phi$  are the solutions to the problem (7). Proposition 2.1 of [6] asserts that  $D_0^{1,2}(\Theta)^\phi$  is compactly embedded in  $L^p(\Theta)$  for  $p \in (2, 2_{n-d}^*)$ , hence, a standard argument shows that the functional  $J_p : D_0^{1,2}(\Theta)^\phi \rightarrow \mathbb{R}$  satisfies the Palais-Smale condition. Therefore,  $J_p$  attains its minimum on  $\mathcal{N}_p^\phi$ . Moreover, as the functional is even and has the mountain pass geometry, the symmetric mountain pass theorem [2] yields the existence of an unbounded sequence of critical values for  $J_p$  in  $D_0^{1,2}(\Theta)^\phi$ .  $\square$

We now derive a multiplicity result for the supercritical problem  $(\wp_p)$ . Assume that the closure of  $\Theta$  is contained in  $\mathbb{R}^{n-1} \times (0, \infty)$  and, for  $m \geq 1$ , let

$$\lambda_1^\phi := \inf_{\substack{u \in D_0^{1,2}(\Theta)^\phi \\ u \neq 0}} \frac{\int_\Theta x_n^m |\nabla u|^2}{\int_\Theta x_n^m u^2}. \quad (9)$$

As the  $n$ -th coordinate  $x_n$  of  $x$  is positive for every  $x \in \overline{\Theta}$ , from the Poincaré inequality we obtain that  $\lambda_1^\phi > 0$ .

**THEOREM 2.3.** *If  $\lambda \in (-\lambda_1^\phi, \infty)$  and  $p \in (2, 2_{n-d}^*)$ , then the problem  $(\wp_p)$  has a least energy solution and an unbounded sequence of solutions in*

$$\Omega := \{(y, z) \in \mathbb{R}^{n-1} \times \mathbb{R}^{m+1} : (y, |z|) \in \Theta\},$$

which satisfy

$$v(\gamma y, \varrho z) = \phi(\gamma)v(y, z) \quad \forall \gamma \in \Gamma, \varrho \in O(m+1), (y, z) \in \Omega. \quad (10)$$

*Proof.* A straightforward computation shows that  $v$  is a solution to the problem  $(\wp_p)$  in  $\Omega$  which satisfies (10) if and only if the function  $u$  given by  $v(y, z) = u(y, |z|)$  is a solution to the problem (7) with  $a(x) := x_n^m =: c(x)$  and  $b(x) := \lambda x_n^m$ . Moreover,  $v$  has minimal energy if and only if  $u$  does. Note that (6) is satisfied if  $\lambda \in (-\lambda_1^\phi, \infty)$ . So this result follows from Theorem 2.2.  $\square$

For  $p \in (2, 2_{n-d}^*)$  let  $u_p$  be a least energy solution to the problem (7). Fix  $q \in (2, 2_{n-d}^*)$  and let  $t_{q,p} \in (0, \infty)$  be such that  $\tilde{u}_p := t_{q,p} u_p \in \mathcal{N}_q^\phi$ , i.e.,

$$t_{q,p} = \left( \frac{\|u_p\|_{a,b}^2}{|u_p|_{c;q}^q} \right)^{\frac{1}{q-2}} = \left( \frac{|u_p|_{c;p}^p}{|u_p|_{c;q}^q} \right)^{\frac{1}{q-2}}. \quad (11)$$

We will show that  $\lim_{p \rightarrow q} J_q(\tilde{u}_p) = \ell_q^\phi$ . The proof is similar to that of Proposition 2.5 in [6]. We give the details for the reader's convenience, starting with the following lemma.

LEMMA 2.4. *If  $p_k, q \in (2, 2_{n-d}^*)$ ,  $p_k \rightarrow q$ , and  $(u_k)$  is a bounded sequence in  $D_0^{1,2}(\Theta)^\phi$ , then*

$$\lim_{k \rightarrow \infty} \int_{\Theta} (c(x) |u_k|^{p_k} - c(x) |u_k|^q) dx = 0.$$

*Proof.* By the mean value theorem, for each  $x \in \Theta$ , there exists  $q_k(x)$  between  $p_k$  and  $q$  such that

$$\left| |u_k(x)|^{p_k} - |u_k(x)|^q \right| = |\ln |u_k(x)|| |u_k(x)|^{q_k(x)} |p_k - q|.$$

Fix  $r > 0$  such that  $[q - r, q + r] \subset (2, 2_{n-d}^*)$ . Then, for some positive constant  $C$  and  $k$  large enough,

$$\left| \ln |u_k| \right| |u_k|^{q_k} \leq \begin{cases} \ln |u_k| |u_k|^{q+r} & \leq C |u_k|^{2_{n-d}^*} & \text{if } |u_k| \geq 1, \\ \left( \ln \frac{1}{|u_k|} \right) |u_k|^{q-r} & \leq C |u_k|^2 & \text{if } |u_k| \leq 1. \end{cases}$$

As  $D_0^{1,2}(\Theta)^\phi$  is continuously embedded in  $L^p(\Theta)$  for  $p \in [2, 2_{n-d}^*]$ , we obtain

$$\begin{aligned} \left| \int_{\Theta} c(|u_k|^{p_k} - |u_k|^q) \right| &\leq |c|_{\infty} \left( \int_{|u_k| \leq 1} ||u_k|^{p_k} - |u_k|^q| + \int_{|u_k| > 1} ||u_k|^{p_k} - |u_k|^q| \right) \\ &\leq |c|_{\infty} C |p_k - q| \int_{\Theta} (|u_k|^2 + |u_k|^{2_{n-d}^*}) \\ &\leq \bar{C} |p_k - q| \|u_k\|^{2_{n-d}^*} \end{aligned}$$

for some positive constant  $\bar{C}$ , where  $|c|_{\infty} := \sup_{x \in \Theta} |c(x)|$ . Since  $(u_k)$  is bounded in  $D_0^{1,2}(\Theta)$ , our claim follows.  $\square$

PROPOSITION 2.5. For  $q \in (2, 2_{n-d}^*)$  we have that

$$\lim_{p \rightarrow q} \ell_p^\phi = \ell_q^\phi, \quad \lim_{p \rightarrow q} t_{q,p} = 1, \quad \lim_{p \rightarrow q} J_q(\tilde{u}_p) = \ell_q^\phi.$$

*Proof.* Set

$$S_p^\phi := \inf_{u \in D_0^{1,2}(\Omega)^\phi \setminus \{0\}} \frac{\|u\|_{a,b}^2}{|u|_{c;p}^2}.$$

It is easy to see that  $\ell_p^\phi = \frac{p-2}{2p} (S_p^\phi)^{\frac{p}{p-2}}$ . So, to prove the first identity, it suffices to show that  $\lim_{p \rightarrow q} S_p^\phi = S_q^\phi$ . From Hölder's inequality we get that  $|u|_{c;q} \leq |c|_1^{(p-q)/pq} |u|_{c;p}$  if  $p > q$ . Hence,  $S_q^\phi \geq |c|_1^{2(q-p)/pq} S_p^\phi$  if  $p > q$ . So, as  $p$  approaches  $q$  from the right, we have that

$$\limsup_{p \rightarrow q^+} S_p^\phi \leq S_q^\phi.$$

Assume that  $\liminf_{p \rightarrow q^+} S_p^\phi < S_q^\phi$ . Then, there exist  $\varepsilon > 0$  and sequences  $(p_k)$  in  $(q, 2_{n-d}^*)$  and  $(u_k)$  in  $D_0^{1,2}(\Omega)^\phi$  with  $|u_k|_{c;p_k} = 1$  such that  $\|u_k\|_{a,b}^2 < S_q^\phi - \varepsilon$ . Lemma 2.4 implies that  $\frac{\|u_k\|_{a,b}^2}{|u_k|_{c;q}^2} < S_q^\phi$  for  $k$  large enough, contradicting the definition of  $S_q^\phi$ . This proves that

$$\lim_{p \rightarrow q^+} S_p^\phi = S_q^\phi.$$

The corresponding statement when  $p$  approaches  $q$  from the left is proved in a similar way. Since  $J_p(u_p) = \frac{p-2}{2p} \|u_p\|_{a,b}^2 = \ell_p^\phi$  we have that  $(u_p)$  is bounded in  $D_0^{1,2}(\Omega)^\phi$  for  $p$  close to  $q$ . Lemma 2.4 applied to (11) yields  $\lim_{p \rightarrow q} t_{q,p} = 1$ . It follows that  $\lim_{p \rightarrow q} J_q(\tilde{u}_p) = \lim_{p \rightarrow q} \frac{q-2}{2q} \|t_{q,p} u_p\|_{a,b}^2 = \ell_q^\phi$ , as claimed.  $\square$

### 3. Minimizing sequences for the critical problem

In this section we analyze the behavior of the minimizing sequences for the problem (7) when  $p$  is the critical exponent  $2_n^* = \frac{2n}{n-2}$ . The solutions to the limit problem (3) will play a crucial role in this analysis. We denote the energy functional associated to (3) by

$$J_\infty(u) := \frac{1}{2} \|u\|^2 - \frac{1}{2_n^*} |u|_{2_n^*}^{2_n^*}$$

and, for any closed subgroup  $K$  of  $\Gamma$ , we set

$$\begin{aligned} D^{1,2}(\mathbb{R}^n)^{\phi|K} &:= \{u \in D^{1,2}(\mathbb{R}^n) : u(\gamma z) = \phi(\gamma)u(z) \ \forall \gamma \in K, z \in \mathbb{R}^n\}, \\ \mathcal{N}_\infty^{\phi|K} &:= \{u \in D^{1,2}(\mathbb{R}^n)^{\phi|K} : u \neq 0, \|u\|^2 = |u|_{2_n^*}^{2_n^*}\}, \\ \ell_\infty^{\phi|K} &:= \inf_{u \in \mathcal{N}_\infty^{\phi|K}} J_\infty(u). \end{aligned}$$

If  $K = \Gamma$  we write  $\mathcal{N}_\infty^\phi$  and  $\ell_\infty^\phi$  instead of  $\mathcal{N}_\infty^{\phi|K}$  and  $\ell_\infty^{\phi|K}$ .

Recall that the  $\Gamma$ -orbit of a point  $x \in \mathbb{R}^n$  is the set  $\Gamma x := \{\gamma x : \gamma \in \Gamma\}$ , and its isotropy group is  $\Gamma_x := \{\gamma \in \Gamma : \gamma x = x\}$ . Then,  $\Gamma x$  is  $\Gamma$ -homeomorphic to the homogeneous space  $\Gamma/\Gamma_x$ . In particular, the cardinality of  $\Gamma x$  is the index of  $\Gamma_x$  in  $\Gamma$ , which is usually denoted by  $|\Gamma/\Gamma_x|$ . If  $\Gamma x = \{x\}$  then  $x$  is said to be a fixed point of  $\Gamma$ . We denote

$$\Theta^\Gamma := \{x \in \Theta : x \text{ is a fixed point of } \Gamma\}.$$

For simplicity, we will write  $J_*$ ,  $\mathcal{N}_*^\phi$  and  $\ell_*^\phi$  instead of  $J_{2_n^*}$ ,  $\mathcal{N}_{2_n^*}^\phi$  and  $\ell_{2_n^*}^\phi$ .

**THEOREM 3.1.** *Let  $(u_k)$  be a sequence in  $\mathcal{N}_*^\phi$  such that  $J_*(u_k) \rightarrow \ell_*^\phi$ . Then, after passing to a subsequence, one of the following two possibilities occurs:*

1.  $(u_k)$  converges strongly in  $D_0^{1,2}(\Theta)$  to a minimizer of  $J_*$  on  $\mathcal{N}_*^\phi$ .
2. There exist a closed subgroup  $K$  of finite index in  $\Gamma$ , a sequence  $(\zeta_k)$  in  $\Theta$ , a sequence  $(\varepsilon_k)$  in  $(0, \infty)$  and a nontrivial solution  $\omega$  to the problem (3) with the following properties:

- (a)  $\Gamma_{\zeta_k} = K$  for all  $k \in \mathbb{N}$ , and  $\zeta_k \rightarrow \zeta$ ,
- (b)  $\varepsilon_k^{-1} \text{dist}(\zeta_k, \partial\Theta) \rightarrow \infty$  and  $\varepsilon_k^{-1} |\alpha\zeta_k - \beta\zeta_k| \rightarrow \infty$  for all  $\alpha, \beta \in \Gamma$  with  $\alpha^{-1}\beta \notin K$ ,
- (c)  $\omega(\gamma z) = \phi(\gamma)\omega(z)$  for all  $\gamma \in K$ ,  $z \in \mathbb{R}^n$ , and  $J_\infty(\omega) = \ell_\infty^{\phi|K}$ ,
- (d)  $\lim_{k \rightarrow \infty} \left\| u_k - \sum_{[\gamma] \in \Gamma/K} \phi(\gamma) \left( \frac{a(\zeta)}{c(\zeta)} \right)^{\frac{n-2}{4}} \varepsilon_k^{\frac{2-n}{2}} (\omega \circ \gamma^{-1}) \left( \frac{\cdot - \gamma\zeta_k}{\varepsilon_k} \right) \right\| = 0$ ,
- (e)  $\ell_*^\phi = \min_{x \in \Theta} \frac{a(x)^{n/2}}{c(x)^{(n-2)/2}} |\Gamma/\Gamma_x| \ell_\infty^{\phi|_{\Gamma_x}} = \frac{a(\zeta)^{n/2}}{c(\zeta)^{(n-2)/2}} |\Gamma/K| J_\infty(\omega)$ .

*Proof.* The proof is exactly the same as that of Theorem 2.5 in [5], omitting the first two lines.  $\square$

Let us state an interesting special case of Theorem 3.1.

**COROLLARY 3.2.** *Assume that every  $\Gamma$ -orbit in  $\Theta$  is either infinite or a fixed point. Let  $(u_k)$  be a sequence in  $\mathcal{N}_*^\phi$  such that  $J_*(u_k) \rightarrow \ell_*^\phi$ . Then, after passing to a subsequence, one of the following statements holds true:*

1.  $(u_k)$  converges strongly in  $D_0^{1,2}(\Theta)$  to a minimizer of  $J_*$  on  $\mathcal{N}_*^\phi$ .

2. There exist a sequence  $(\zeta_k)$  in  $\Theta^\Gamma$ , a sequence  $(\varepsilon_k)$  in  $(0, \infty)$  and a nontrivial  $\phi$ -equivariant solution  $\omega$  to the limit problem (3) such that  $\zeta_k \rightarrow \zeta \in \Theta$ ,  $\varepsilon_k^{-1} \text{dist}(\zeta_k, \partial\Theta) \rightarrow \infty$ ,  $J_\infty(\omega) = \ell_\infty^\phi$ ,

$$\lim_{k \rightarrow \infty} \left\| u_k - \left( \frac{a(\zeta)}{c(\zeta)} \right)^{\frac{n-2}{4}} \varepsilon_k^{\frac{2-n}{2}} \omega \left( \frac{\cdot - \zeta_k}{\varepsilon_k} \right) \right\| = 0,$$

and

$$\frac{a(\zeta)^{n/2}}{c(\zeta)^{(n-2)/2}} = \min_{x \in \Theta^\Gamma} \frac{a(x)^{n/2}}{c(x)^{(n-2)/2}}.$$

In particular, if every  $\Gamma$ -orbit in  $\Theta$  has positive dimension, then (1) must hold true.

*Proof.* Since the group  $K = \Gamma_{\zeta_k}$ , given by case 2 of Theorem 3.1, has finite index in  $\Gamma$  and this index is the cardinality of the  $\Gamma$ -orbit of  $\zeta_k$ , our assumption implies that  $K = \Gamma$  and  $\zeta_k$  is a fixed point. So case 2 of Theorem 3.1 reduces to case 2 of this corollary.  $\square$

Note that the functions  $a$  and  $c$  determine the location of the concentration point  $\zeta$ .

It was shown in [4, Theorem 2.3] that, if  $a = c = 1$ ,  $b = 0$  and  $\Theta^\Gamma \neq \emptyset$ , then  $\ell_*^\phi$  is not attained by  $J_*$  on  $\mathcal{N}_*^\phi$ . So, if every  $\Gamma$ -orbit in  $\Theta \setminus \Theta^\Gamma$  has positive dimension, statement 2 of Corollary 3.2 must hold true.

In the following section we will state a nonexistence result which allows us to obtain information on the limit profile of solutions to the problem  $(\wp_p)$ .

#### 4. Blow-up at the higher critical exponents

Throughout this section we will assume that  $\Theta$  is a  $\Gamma$ -invariant bounded smooth domain in  $\mathbb{R}^n$  whose closure is contained in  $\mathbb{R}^{n-1} \times (0, \infty)$ . Then, the points in  $\{0\} \times (0, \infty)$  must be fixed points of  $\Gamma$ , so  $\mathbb{R}^{n-1} \times \{0\}$  is  $\Gamma$ -invariant and we may regard  $\Gamma$  as a subgroup of  $O(n-1)$ . We will also assume that  $\Theta \setminus \Theta^\Gamma$  and  $\Theta^\Gamma$  are nonempty, and that every  $\Gamma$ -orbit in  $\Theta \setminus \Theta^\Gamma$  has positive dimension. As before,  $\phi : \Gamma \rightarrow \mathbb{Z}_2$  will be a continuous homomorphism which satisfies assumption (5).

We set

$$\Theta_\delta := \{x \in \Theta : \text{dist}(x, \Theta^\Gamma) > \delta\} \text{ if } \delta > 0, \quad \text{and} \quad \Theta_0 := \Theta,$$

and we fix  $\delta_0 > 0$  such that  $\Theta_{\delta_0} \neq \emptyset$ . For  $m \geq 1$  and  $\delta \in [0, \delta_0)$ , we consider the problem

$$(\wp_{\delta,p}^\#) \quad \begin{cases} -\text{div}(x_n^m \nabla u) + \lambda x_n^m u = x_n^m |u|^{p-2} u & \text{in } \Theta_\delta, \\ u = 0 & \text{on } \partial\Theta_\delta, \\ u(\gamma x) = \phi(\gamma)u(x), & \forall \gamma \in \Gamma, \ x \in \Theta_\delta, \end{cases}$$

where  $x_n^m$  denotes the function  $x = (x_1, \dots, x_n) \mapsto x_n^m$ , and  $\lambda \in (-\lambda_1^\phi, \infty)$ , with  $\lambda_1^\phi$  as defined in (9). Then, the operator  $-\operatorname{div}(x_n^m \nabla) + \lambda x_n^m$  is coercive in  $D_0^{1,2}(\Theta)^\phi$ . So the data of this problem satisfy all assumptions stated at the beginning of Section 2.

Theorem 2.2 asserts that the problem  $(\varphi_{\delta,p}^\#)$  has a least energy solution  $u_{\delta,p}$  if  $\delta \in (0, \delta_0)$  and  $p \in (2, 2_{n-\vartheta}^*)$ , where

$$\vartheta := \min\{\dim(\Gamma x) : x \in \Theta \setminus \Theta^\Gamma\}.$$

Note that, by assumption,  $\vartheta > 0$ . On the other hand, for  $\delta = 0$  and  $p = 2_n^*$ , the following nonexistence result was proved in [5].

**THEOREM 4.1.** *If  $\operatorname{dist}(\Theta^\Gamma, \mathbb{R}^{n-1} \times \{0\}) = \operatorname{dist}(\Theta, \mathbb{R}^{n-1} \times \{0\})$ , then there exists  $\lambda_* \in (-\lambda_1^\phi, 0]$  such that, if  $\lambda \in (\lambda_*, \infty) \cup \{0\}$ , the critical problem  $(\varphi_{0,2_n^*}^\#)$  does not have a least energy solution.*

*Moreover,  $\lambda_* < 0$  if  $m \geq 2$ .*

*Proof.* See Theorem 3.2 in [5].  $\square$

For  $\delta \in (0, \delta_0)$  and  $p \in (2, 2_{n-\vartheta}^*)$ , let  $J_{\delta,p} : D_0^{1,2}(\Theta_\delta)^\phi \rightarrow \mathbb{R}$  be the variational functional and  $\mathcal{N}_{\delta,p}^\phi$  be the Nehari manifold associated to the problem  $(\varphi_{\delta,p}^\#)$ , and set

$$\ell_{\delta,p}^\phi := \inf\{J_{\delta,p}(u) : u \in \mathcal{N}_{\delta,p}^\phi\}.$$

We write  $J_*$ ,  $\mathcal{N}_*^\phi$  and  $\ell_*^\phi$  for the variational functional, the Nehari manifold and the infimum associated to the critical problem  $(\varphi_{0,2_n^*}^\#)$  in the whole domain  $\Theta$ . Extending each function in  $\mathcal{N}_{\delta,2_n^*}^\phi$  by 0 outside of  $\Theta_\delta$ , we have that  $\mathcal{N}_{\delta,2_n^*}^\phi \subset \mathcal{N}_*^\phi$  and  $J_{\delta,2_n^*}(u) = J_*(u)$  for every  $u \in \mathcal{N}_{\delta,2_n^*}^\phi$ . Hence,  $\ell_*^\phi \leq \ell_{\delta,2_n^*}^\phi$ .

**LEMMA 4.2.**  $\ell_{\delta,2_n^*}^\phi \rightarrow \ell_*^\phi$  as  $\delta \rightarrow 0$ .

*Proof.* Let  $X := (\mathbb{R}^n)^\Gamma$  and  $Y$  be its orthogonal complement in  $\mathbb{R}^n$ . Since  $\Theta \setminus \Theta^\Gamma \neq \emptyset$  and every  $\Gamma$ -orbit in  $\Theta \setminus \Theta^\Gamma$  has positive dimension, we have that  $\dim(Y) \geq 2$ .

We claim that there are radial functions  $\chi_k \in \mathcal{C}_c^\infty(Y)$  such that  $\chi_k(y) = 1$  if  $|y| \leq \frac{1}{k}$ ,

$$\lim_{k \rightarrow \infty} \int_Y |\chi_k|^2 = 0 \quad \text{and} \quad \lim_{k \rightarrow \infty} \int_Y |\nabla \chi_k|^2 = 0. \quad (12)$$

To show this, we choose a radial function  $g \in \mathcal{C}_c^\infty(Y)$  such that  $g(y) = 1$  if  $|y| \leq 1$  and  $g(y) = 0$  if  $|y| \geq 2$ , and we set  $g_k(y) := g(ky)$ . Define

$$\chi_k(y) := \frac{1}{\sigma_k} \sum_{j=1}^k \frac{g_j(y)}{j}, \quad \text{where } \sigma_k := \sum_{j=1}^k \frac{1}{j}.$$

Clearly,  $\chi_k(y) = 1$  if  $|y| \leq \frac{1}{k}$  and  $\chi_k(y) = 0$  if  $|y| \geq 2$ . As  $\dim(Y) \geq 2$ , we have that  $\int_Y |\nabla g_k|^2 \leq \int_Y |\nabla g|^2$ . Hence, for some positive constant  $C$ ,

$$\int_Y |\nabla \chi_k|^2 \leq \frac{C}{\sigma_k^2} \sum_{j=1}^k \frac{1}{j^2} \rightarrow 0 \quad \text{as } k \rightarrow \infty.$$

Finally, as all functions  $\chi_k$  are supported in the closed ball of radius 2 in  $Y$ , the Poincaré inequality yields

$$\int_Y |\chi_k|^2 \leq C \int_Y |\nabla \chi_k|^2 \rightarrow 0,$$

and our claim is proved.

Given  $\varepsilon > 0$  we choose  $\psi \in \mathcal{N}_*^\phi$  such that  $J_*(\psi) < \ell_*^\phi + \frac{\varepsilon}{2}$ . For  $(x, y) \in X \times Y$ , we define  $\psi_k(x, y) := (1 - \chi_k(y))\psi(x, y)$ . Note that, as  $\chi_k$  is radial and  $\psi$  is  $\phi$ -equivariant,  $\psi_k$  is also  $\phi$ -equivariant. Moreover, the identities (12) easily imply that  $\psi_k \rightarrow \psi$  in  $D_0^{1,2}(\Theta)$ . So, for  $k$  large enough, there exists  $t_k \in (0, \infty)$  such that  $\tilde{\psi}_k := t_k \psi_k \in \mathcal{N}_*^\phi$  and  $t_k \rightarrow 1$ . Hence,  $\tilde{\psi}_k \rightarrow \psi$  in  $D_0^{1,2}(\Theta)$ , and we may choose  $k_0 \in \mathbb{N}$  such that  $J_*(\tilde{\psi}_{k_0}) < \ell_*^\phi + \varepsilon$ . Observe that  $\text{supp}(\tilde{\psi}_k) = \text{supp}(\psi_k) \subset \Theta_\delta$  if  $\delta < \frac{1}{k}$ . So  $\tilde{\psi}_k \in \mathcal{N}_{\delta, 2_n^*}^\phi$  if  $\delta < \frac{1}{k}$ . It follows that

$$\ell_*^\phi \leq \ell_{\delta, 2_n^*}^\phi \leq J_{\delta, 2_n^*}(\tilde{\psi}_{k_0}) = J_*(\tilde{\psi}_{k_0}) < \ell_*^\phi + \varepsilon \quad \forall \delta \in \left(0, \frac{1}{k_0}\right).$$

This finishes the proof.  $\square$

Set  $N := n + m$  and

$$\Omega_\delta := \{(y, z) \in \mathbb{R}^{n-1} \times \mathbb{R}^{m+1} : (y, |z|) \in \Theta_\delta\}, \quad \delta \in [0, \delta_0].$$

Note that  $\Omega_\delta$  is  $[\Gamma \times O(m+1)]$ -invariant, i.e.,  $(\gamma y, \varrho z) \in \Omega_\delta$  for every  $(y, z) \in \Omega_\delta$ ,  $\gamma \in \Gamma$ ,  $\varrho \in O(m+1)$ . A straightforward computation shows that  $u_{\delta, p}$  is a least energy solution to the problem  $(\varphi_{\delta, p}^\#)$  if and only if  $v_{\delta, p}(y, z) := u_{\delta, p}(y, |z|)$  is a least energy solution to the problem

$$(\varphi_{\delta, p}) \quad \begin{cases} -\Delta v + \lambda v = |v|^{p-2}v & \text{in } \Omega_\delta, \\ v = 0 & \text{on } \partial\Omega_\delta, \\ v(\gamma y, \varrho z) = \phi(\gamma)v(y, z), \quad \forall \gamma \in \Gamma, \varrho \in O(m+1), (y, z) \in \Omega_\delta. \end{cases}$$

Therefore, for every  $\lambda \in (-\lambda_1^\phi, \infty)$ ,  $\delta \in (0, \delta_0)$  and  $p \in (2, 2_{n-d}^*)$ , the problem  $(\varphi_{\delta, p})$  has a least energy solution. The following results describe its limit profile.

**THEOREM 4.3.** *For  $\delta \in (0, \delta_0)$  let  $v_{\delta,*}$  be a least energy solution to the problem  $(\wp_{\delta, 2_{N,m}^*})$ . Assume that*

$$\text{dist}(\Theta^\Gamma, \mathbb{R}^{n-1} \times \{0\}) = \text{dist}(\Theta, \mathbb{R}^{n-1} \times \{0\}).$$

*Then, there exists  $\lambda_* \leq 0$  such that, if  $\lambda \in (\lambda_*, \infty) \cup \{0\}$ , there exist sequences  $(\delta_k)$  in  $(0, \delta_0)$ ,  $(\varepsilon_k)$  in  $(0, \infty)$ ,  $(\zeta_k)$  in  $\Theta^\Gamma$ , and a nontrivial solution  $\omega$  to the limit problem (3) such that*

(i)  $\delta_k \rightarrow 0$ ,  $\varepsilon_k^{-1} \text{dist}(\zeta_k, \partial\Theta) \rightarrow \infty$ , and  $\zeta_k \rightarrow \zeta$  with

$$\text{dist}(\zeta, \mathbb{R}^{n-1} \times \{0\}) = \text{dist}(\Theta, \mathbb{R}^{n-1} \times \{0\}),$$

(ii)  $\omega$  is  $\phi$ -equivariant and has minimal energy among all nontrivial  $\phi$ -equivariant solutions to the problem (3),

(iii)  $v_{\delta_k,*} = \tilde{\omega}_{\varepsilon_k, \zeta_k} + o(1)$  in  $D^{1,2}(\mathbb{R}^N)$ , where

$$\tilde{\omega}_{\varepsilon_k, \zeta_k}(y, z) := \varepsilon_k^{(2-n)/2} \omega \left( \frac{(y, |z|) - \zeta_k}{\varepsilon_k} \right).$$

Moreover,  $\lambda_* < 0$  if  $m \geq 2$ .

*Proof.* Let  $\lambda_*$  be the number given by Theorem 4.1. Fix  $\lambda \in (\lambda_*, \infty) \cup \{0\}$ , and let  $u_{\delta,*}$  be the least energy solution to the problem  $(\wp_{\delta, 2_n^*}^\#)$  given by  $v_{\delta,*}(y, z) = u_{\delta,*}(y, |z|)$ . Choose a sequence  $\delta_k \rightarrow 0$  and set  $u_k := u_{\delta_k,*}$ . Then,  $u_k \in \mathcal{N}_*^\phi$  and, by Lemma 4.2,  $J_*(u_k) \rightarrow \ell_*^\phi$ . It follows from Corollary 3.2 and Theorem 4.1 that, after passing to a subsequence, there exist sequences  $(\varepsilon_k)$  in  $(0, \infty)$  and  $(\zeta_k)$  in  $\Theta^\Gamma$ , and a nontrivial  $\phi$ -equivariant solution  $\omega$  to the limit problem (3) such that  $\zeta_k \rightarrow \zeta$ ,  $\varepsilon_k^{-1} \text{dist}(\zeta_k, \partial\Theta) \rightarrow \infty$ ,  $J_\infty(\omega) = \ell_\infty^\phi$ ,

$$\lim_{k \rightarrow \infty} \left\| u_k - \varepsilon_k^{\frac{2-n}{2}} \omega \left( \frac{\cdot - \zeta_k}{\varepsilon_k} \right) \right\| = 0, \quad (13)$$

and

$$[\text{dist}(\zeta, \mathbb{R}^{n-1} \times \{0\})] = \min_{x \in \overline{\Theta}} [\text{dist}(x, \mathbb{R}^{n-1} \times \{0\})].$$

Equation (13) implies that  $v_{\delta_k,*}$  satisfies (3). This concludes the proof.  $\square$

**THEOREM 4.4.** *For  $\delta \in (0, \delta_0)$  and  $p \in (2_{N,m}^*, 2_{N,m+d}^*)$  let  $v_{\delta,p}$  be a least energy solution to the problem  $(\wp_{\delta,p})$ . Assume that*

$$\text{dist}(\Theta^\Gamma, \mathbb{R}^{n-1} \times \{0\}) = \text{dist}(\Theta, \mathbb{R}^{n-1} \times \{0\}).$$

*Then, there exists  $\lambda_* \leq 0$  such that, if  $\lambda \in (\lambda_*, \infty) \cup \{0\}$ , there exist sequences  $(\delta_k)$  in  $(0, \delta_0)$ ,  $(\varepsilon_k)$  in  $(0, \infty)$ ,  $(p_k)$  in  $(2_{N,m}^*, 2_{N,m+d}^*)$ , and  $(\zeta_k)$  in  $\Theta^\Gamma$ , and a nontrivial solution  $\omega$  to the limit problem (3) such that*



(i)  $\delta_k \rightarrow 0$ ,  $p_k \rightarrow 2_{N,m}^*$ ,  $\varepsilon_k^{-1} \text{dist}(\zeta_k, \partial\Theta) \rightarrow \infty$ , and  $\zeta_k \rightarrow \zeta$  with

$$\text{dist}(\zeta, \mathbb{R}^{n-1} \times \{0\}) = \text{dist}(\Theta, \mathbb{R}^{n-1} \times \{0\}),$$

(ii)  $\omega$  is  $\phi$ -equivariant and has minimal energy among all nontrivial  $\phi$ -equivariant solutions to the problem (3),

(iii)  $v_{\delta_k, p_k} = \tilde{\omega}_{\varepsilon_k, \zeta_k} + o(1)$  in  $D^{1,2}(\mathbb{R}^N)$ , where

$$\tilde{\omega}_{\varepsilon_k, \zeta_k}(y, z) := \varepsilon_k^{(2-n)/2} \omega \left( \frac{(y, |z|) - \zeta_k}{\varepsilon_k} \right).$$

Moreover,  $\lambda_* < 0$  if  $m \geq 2$ .

*Proof.* Let  $\lambda_*$  be the number given by Theorem 4.1. Fix  $\lambda \in (\lambda_*, \infty) \cup \{0\}$ . Let  $u_{\delta, p}$  be the least energy solution to the problem  $(\varphi_{\delta, p}^\#)$  given by  $v_{\delta, p}(y, z) = u_{\delta, p}(y, |z|)$  and let  $t_{\delta, p} \in (0, \infty)$  be such that  $\tilde{u}_{\delta, p} := t_{\delta, p} u_{\delta, p} \in \mathcal{N}_{\delta, 2_n^*}^\phi \subset \mathcal{N}_*^\phi$ . Proposition 2.5 and Lemma 4.2 allow us to choose  $\delta_k \in (0, \delta_0)$  and  $p_k \in (2_n^*, 2_{n-2}^*)$  such that  $\delta_k \rightarrow 0$ ,  $p_k \rightarrow 2_n^*$ , and  $J_*(\tilde{u}_k) \rightarrow \ell_*^\phi$ , where  $\tilde{u}_k := \tilde{u}_{\delta_k, p_k}$ . The rest of the proof is the same as that of Theorem 4.3  $\square$

Finally, we derive Theorems 1.1 and 1.2 from Theorems 2.3 and 4.4.

*Proof of Theorem 1.1.* Let  $\Gamma := O(n-1)$  and  $\phi$  be the trivial homomorphism  $\phi \equiv 1$ . Then,  $B^\Gamma = B \cap [\{0\} \times (0, \infty)]$ . A  $\phi$ -equivariant function is simply a  $\Gamma$ -invariant function and, as the standard bubble is radial, it is the least energy  $\Gamma$ -invariant solution to the problem (3), which is unique up to translations and dilations. Since  $\dim(\Gamma x) = n-2 \geq 1$  for every  $x \in B \setminus B^\Gamma$ , applying Theorems 2.3 and 4.4 to  $\Theta := B$  with this group action we obtain Theorem 1.1.  $\square$

*Proof of Theorem 1.2.* For  $n \geq 5$ , let  $\Gamma$  be the subgroup of  $O(n-1)$  generated by  $\{e^{i\vartheta}, \alpha, \tau : \vartheta \in [0, 2\pi), \alpha \in O(n-5)\}$  acting on a point  $y = (\eta, \xi) \in \mathbb{C}^2 \times \mathbb{R}^{n-5} \cong \mathbb{R}^{n-1}$ ,  $\eta = (\eta_1, \eta_2) \in \mathbb{C} \times \mathbb{C}$ , as

$$e^{i\vartheta} y := (e^{i\vartheta} \eta, \xi), \quad \alpha y := (\eta, \alpha \xi), \quad \tau y := (-\bar{\eta}_2, \bar{\eta}_1, \xi),$$

and let  $\phi$  be the homomorphism given by  $\phi(e^{i\vartheta}) = 1 = \phi(\alpha)$  and  $\phi(\tau) = -1$ . Then,  $B^\Gamma = B \cap [\{0\} \times (0, \infty)]$ . If  $n = 5$  then  $\dim(\Gamma y) = 1$  for every  $y \in \mathbb{R}^{n-1} \setminus \{0\}$ , whereas for  $n \geq 6$  we have that

$$\dim(\Gamma y) = \begin{cases} n-5 & \text{if } \eta \neq 0 \text{ and } \xi \neq 0, \\ 1 & \text{if } \xi = 0, \\ n-6 & \text{if } \eta = 0. \end{cases}$$

Therefore, if  $n = 5$  or  $n \geq 7$ , we have that  $\dim(\Gamma x) \geq 1$  for every  $x \in B \setminus B^\Gamma$ . Notice that any point  $x_0 = (\eta, \xi) \in B$  with  $\eta \neq 0$  satisfies condition (5). Hence, Theorem 1.2 follows from Theorems 2.3 and 4.4.  $\square$

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## REFERENCES

- [1] N. ACKERMANN, M. CLAPP, AND A. PISTOIA, *Boundary clustered layers near the higher critical exponents*, J. Differential Equations **254** (2013), no. 10, 4168–4193.
- [2] A. AMBROSETTI AND P. H. RABINOWITZ, *Dual variational methods in critical point theory and applications*, J. Functional Analysis **14** (1973), 349–381.
- [3] J. BRACHO, M. CLAPP, AND W. MARZANTOWICZ, *Symmetry breaking solutions of nonlinear elliptic systems*, Topol. Methods Nonlinear Anal. **26** (2005), no. 1, 189–201.
- [4] M. CLAPP, *Entire nodal solutions to the pure critical exponent problem arising from concentration*, J. Differential Equations **261** (2016), no. 6, 3042–3060.
- [5] M. CLAPP AND J. FAYA, *Concentration with a single sign changing layer at the higher critical exponents*, Advances in Nonlinear Analysis (2016), DOI 10.1515/anona-2016-0056.
- [6] M. CLAPP AND F. PACELLA, *Existence and asymptotic profile of nodal solutions to supercritical problems*, Advanced Nonlinear Studies (2016), DOI 10.1515/ans-2016-6009.
- [7] M. CLAPP AND A. PISTOIA, *Symmetries, Hopf fibrations and supercritical elliptic problems*, Contemp. Math. (Amer. Math. Soc.) **656** (2016), 1–12.
- [8] M. DEL PINO, P. FELMER, AND M. MUSSO, *Multi-bubble solutions for slightly super-critical elliptic problems in domains with symmetries*, Bull. London Math. Soc. **35** (2003), no. 4, 513–521.
- [9] M. DEL PINO, P. FELMER, AND M. MUSSO, *Two-bubble solutions in the super-critical Bahri-Coron’s problem*, Calc. Var. Partial Differential Equations **16** (2003), no. 2, 113–145.
- [10] M. DEL PINO, M. MUSSO, AND F. PACARD, *Bubbling along boundary geodesics near the second critical exponent*, J. Eur. Math. Soc. (JEMS) **12** (2010), no. 6, 1553–1605.
- [11] M. DEL PINO, M. MUSSO, F. PACARD, AND A. PISTOIA, *Large energy entire solutions for the Yamabe equation*, J. Differential Equations **251** (2011), no. 9, 2568–2597.
- [12] E. HEBEY AND M. VAUGON, *Sobolev spaces in the presence of symmetries*, J. Math. Pures Appl. (9) **76** (1997), no. 10, 859–881.
- [13] S. KIM AND A. PISTOIA, *Boundary towers of layers for some supercritical problems*, J. Differential Equations **255** (2013), no. 8, 2302–2339.
- [14] R. MOLLE AND D. PASSASEO, *Positive solutions for slightly super-critical elliptic equations in contractible domains*, C. R. Math. Acad. Sci. Paris **335** (2002), no. 5, 459–462.

- [15] R. MOLLE AND D. PASSASEO, *Positive solutions of slightly supercritical elliptic equations in symmetric domains*, Ann. Inst. H. Poincaré Anal. Non Linéaire **21** (2004), no. 5, 639–656.
- [16] M. MUSSO AND A. PISTOIA, *Persistence of Coron's solution in nearly critical problems*, Ann. Sc. Norm. Super. Pisa Cl. Sci. (5) **6** (2007), no. 2, 331–357.
- [17] M. MUSSO AND J. WEI, *Sign-changing blowing-up solutions for supercritical Bahri-Coron's problem*, Calc. Var. Partial Differential Equations **55** (2016), no. 1, 1–39.
- [18] D. PASSASEO, *Nonexistence results for elliptic problems with supercritical nonlinearity in nontrivial domains*, J. Funct. Anal. **114** (1993), no. 1, 97–105.
- [19] D. PASSASEO, *New nonexistence results for elliptic equations with supercritical nonlinearity*, Differential Integral Equations **8** (1995), no. 3, 577–586.
- [20] A. PISTOIA AND O. REY, *Multiplicity of solutions to the supercritical Bahri-Coron's problem in pierced domains*, Adv. Differential Equations **11** (2006), no. 6, 647–666.

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